

# Anisotropic electronic structure in quasi-one-dimensional $K_{0.3}MoO_3$ : An angle-dependent x-ray absorption study

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The electronic structure of quasi-one-dimensional (quasi-1D) blue bronze,  $K_{0.3}MoO_3$ , was investigated by angle-dependent x-ray absorption near-edge structure (XANES) spectroscopy at O and K  $K$  and Mo  $L_3$  edges along the quasi-1D  $MoO_6$  octahedron-chain direction, i.e., the  $b$  axis, and the octahedron-in-plane direction, i.e., the  $d$  axis, well below its Peierls phase transition temperature (180 K). The O  $K$ -edge XANES spectra indicate that the angle dependence of O  $2p$ -Mo  $4d$  hybridization, especially those with the  $\pi^*$  character, is more significant along the  $b$  axis than along the  $d$  axis. Similar trend is also observed in the Mo  $L_3$ -edge XANES spectra. The K  $K$ -edge XANES spectra reveal anisotropic effect of hybridization of K  $4p$  states with O  $2p$  states on the  $MoO_6$  octahedron. © 2007 American Institute of Physics. [DOI: 10.1063/1.2756358]

Molybdenum bronzes belong to a class of solid oxides with intense color and metallic luster.<sup>1</sup> Of these,  $K_{0.3}MoO_3$  (KMO) is a prototype blue bronze that exhibits quasi-one-dimensional (quasi-1D) electron transport property and a metal-semiconductor transition to a charge-density-wave (CDW) phase at 180 K.<sup>1-4</sup> Electrical, magnetic, and optical properties are strongly anisotropic with quasi-1D behavior. Hence, this material has been investigated using various experimental approaches to understand the origin of these properties. Its electrical conductivity is an order of magnitude larger along the chain ( $b$  axis) than along the octahedron-in-plane ( $d$  axis) direction and it is very sensitive to the thermodynamical parameters.<sup>1-4</sup> A detailed understanding of the direction-dependent electronic structure is critical to elucidate the anisotropic behavior of KMO. Such a study will also provide information on the electronic states of O and Mo atoms and the role played by K atoms in its anisotropic electronic properties, which has not been revealed in previous experimental and theoretical investigations.<sup>1-6</sup>

KMO has a centered monoclinic structure and contains  $MoO_6$  chains with K atoms residing between adjacent chains. KMO has 20 f.u. in a unit cell, which has ten distorted  $MoO_6$  octahedra that share edges or corners in a cluster (see Fig. 1)<sup>1,3,7,8</sup> These clusters are connected through three nonequivalent Mo sites in the KMO structure. The Mo(1) octahedron shares its edges and corners with only the Mo(2) and Mo(3) octahedra of the same unit cell. These clusters are linked together by sharing corners along and perpendicular to the  $b$  axis and form into chains. There are two nonequivalent K sites. The vector  $d=2c+a$ , where  $c$  and  $a$  are the lattice parameters, is perpendicular to the  $b$  axis. The unit cell in blue bronze is spanned by vectors  $b$  and  $d$ .<sup>7,8</sup>

When spectroscopic measurements are made along these two axes, the anisotropy can be identified. X-ray absorption near-edge structure (XANES) is an element specific probe that is sensitive to local structure around the absorbing atom and to the coordination geometry of its near neighbors. Linear polarization-dependent XANES can be employed to probe selectively the symmetry of the final states relative to the various crystallographic directions. Duda *et al.*<sup>7</sup> and Sing *et al.*<sup>8</sup> adopted this approach to understand the electronic structure of KMO. Huang *et al.* investigated this compound using the same technique, by varying the temperature and the applied voltage, to elucidate O  $2p$ -Mo  $4d$  hybridized states and the effect of the local atomic structure associated with the dynamic CDW of KMO.<sup>9</sup> This work uses angle-dependent XANES spectroscopy to probe the electronic structures at O, Mo, and K sites, along the  $b$  and  $d$  axes, respectively, by the photoexcitation of electrons from discrete core levels to states above the Fermi level. These mea-

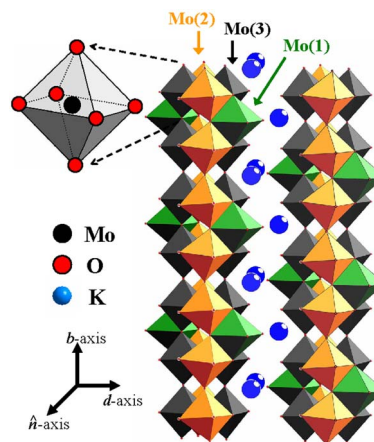


FIG. 1. (Color online) Crystal structure of the prototype blue bronze,  $K_{0.3}MoO_3$ ; nonequivalent Mo sites are denoted by Mo(1), Mo(2), and Mo(3). The zigzag chain structure of K ions is also evident.

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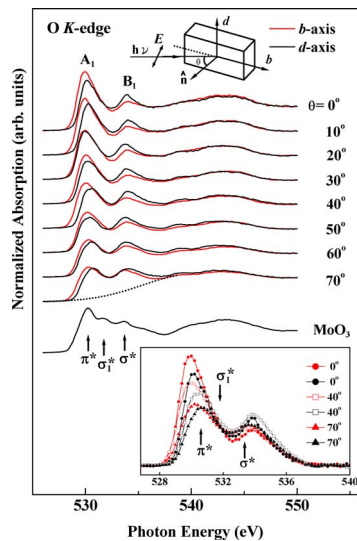


FIG. 2. (Color online) Normalized XANES spectra at the O  $K$  edge measured at various angles and along the  $b$  and  $d$  axes obtained at 140 K. The top inset presents the geometry of measurement and the bottom inset magnifies the variation in spectral features at various angles,  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$ , and along the  $b$  and  $d$  axes.

measurements provide a site-specific probe of the symmetry of unoccupied states and yield information on the local environments at O, Mo, and K sites in the CDW phase of the KMO.

The angle-dependent XANES spectra were obtained at the National Synchrotron Radiation Research Center, Hsinchu, Taiwan. The KMO single crystal was grown using a temperature gradient flux technique and its orientation was determined by x-ray diffraction.<sup>10</sup> A rectangular piece of crystal with an area of  $\sim 2 \times 3$  mm<sup>2</sup> was prepared and the longer side of the rectangle was aligned with the  $b$  axis of the crystal. All these measurements were carried out at 140 K with a temperature stability of  $\pm 1$  K using a continuous flow-type helium cryostat; this temperature was well below the Peierls phase transition temperature of 180 K of KMO.

Figure 2 displays the normalized O  $K$ -edge XANES spectra of KMO and reference MoO<sub>3</sub> along both the  $b$  and  $d$  axes at angles of  $0^\circ$ – $70^\circ$ . The upper part of the inset in Fig. 2 presents the experimental geometry used in this work, in which  $E$  is the polarization of the incoming photons and  $\theta$  is the incidence angle. These spectral features are consistent with the reports available in the literature.<sup>7–9</sup> The O  $K$ -edge XANES spectra reflect transitions from the O  $1s$  core state to the unoccupied O  $2p$  derived states and the states of neighboring atoms, which have  $p$ -symmetry components projected onto the O sites. The positions of the spectral features agree with those of the calculated partial densities of states of KMO using the first-principles method.<sup>6</sup> The two prominent features ( $A_1$  and  $B_1$ ) near 530 and 534 eV are known to correspond to the unoccupied  $\pi^*$  and  $\sigma^*$  bands formed by hybridization between O  $2p_\pi$  and Mo  $4dt_{2g}$  states and between O  $2p_\sigma$  and Mo  $4de_g$  states,<sup>7–9</sup> respectively. The shoulder ( $\sigma_1^*$ ) at  $\sim 532$  eV from the reference compound, MoO<sub>3</sub>, is associated with the anisotropy of the Mo  $4dt_{2g}$  band due to the presence of nonequivalent oxygen atoms in the first coordination shell,<sup>11,12</sup> which is not clearly observed in the spectra of KMO. This result reveals that the environments of O ions differ from those in MoO<sub>3</sub>. The hybridization of the K  $4sp$ –O  $2p$  states reduces the number of unoccupied states of anisotropic O  $2p$ – $4dt_{2g}$  bands centered at  $\sim 532$  eV in the

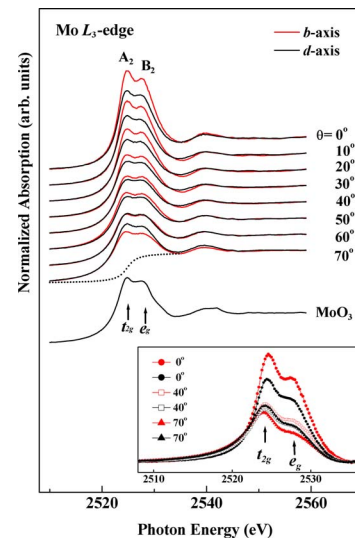


FIG. 3. (Color online) Normalized XANES spectra at the Mo  $L_3$  edge measured at various angles and along the  $b$  and  $d$  axes obtained at 140 K. The inset magnifies the variation in the  $4d$   $t_{2g}$  and  $e_g$  bands at various angles,  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$ , and along the  $b$  and  $d$  axes.

KMO spectra and broadens feature  $B_1$ . The broad humplike feature from  $\sim 537$  to 546 eV was attributed to Mo  $5sp$ –O  $2p$  hybridized states.<sup>7,8,11</sup> The lower inset in Fig. 2 presents magnified near-edge features in the  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$  spectra along the  $b$  and  $d$  axes after subtraction of the Gaussian background indicated by the dashed line in Fig. 2. As revealed in the inset, the intensity of the  $\pi^*$  feature is highest at  $\theta=0^\circ$  and decreases to a minimum at  $\theta=70^\circ$  as the angle increases; the intensity of the  $\sigma^*$  feature varies oppositely. In  $d$ -axis spectra,  $\pi^*$  features slightly shift to higher energies and have a similar trend as those in the  $b$  axis spectra. However, the intensity of the  $\sigma^*$  feature does not show a clear trend. The difference intensities between  $b$ - and  $d$ -axis spectra are substantially larger for smaller angles than those for larger angles, because at small angles the  $E$  field of photons are predominantly parallel with the  $b$  axis and along the  $b$  axis O  $K$ -edge measurement probes a larger number of Mo–O hybridized states due to a larger number of Mo–O bonds along the chain, while along the  $d$  axis the MoO<sub>6</sub> clusters are separated by K ions and consequently O  $K$ -edge measurement probes a smaller number of Mo–O bonds and Mo–O hybridized states.

Figure 3 displays normalized Mo  $L_3$ -edge XANES spectra of KMO and reference MoO<sub>3</sub> at various angles along the  $b$  and  $d$  axes. Spectral features  $A_2$  and  $B_2$  reflect transitions from Mo  $2p_{3/2}$  to  $4d$  orbitals and provide information about unoccupied Mo  $4d$ -derived states.<sup>13</sup> As shown in the figure, the white lines clearly split into  $4dt_{2g}$  and  $4de_g$  bands. The difference between the spectral features of KMO and MoO<sub>3</sub> indicates that MoO<sub>6</sub> octahedra are distorted by the insertion of K atoms between MoO<sub>6</sub> chains, which broadens both  $t_{2g}$  and  $e_g$  bands. The lower inset in Fig. 3 presents magnified near-edge features in the spectra of  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$  along the  $b$  and  $d$  axes, in which the intensities of the white line feature were subtracted with a fitted arctangent function of the continuum step centered at the position of the maximum height, as indicated by the dashed line in Fig. 3. The intensities of the features for the  $b$ - and  $d$ -axis cases decrease with the increase of  $\theta$ . The rate of change of the intensity with respect to  $\theta$  is much larger for the  $b$  axis case than for the

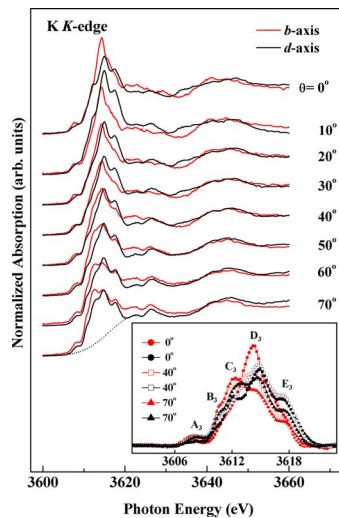


FIG. 4. (Color online) Normalized XANES spectra at the K  $K$  edge measured at various angles and along the  $b$  and  $d$  axes obtained at 140 K. The inset magnifies the variation in spectral features at various angles,  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$ , and along the  $b$  and  $d$  axes.

$d$ -axis case such that at  $\theta=0^\circ$  the intensity for the  $b$ -axis case is the largest, while at  $\theta=70^\circ$  it is the smallest. This trend is consistent with that of the near-edge features at the O  $K$  edge and is originated from the same anisotropic structural property.

The  $\text{MoO}_6$  octahedra and the double  $\text{ReO}_3$ -type chains are common to  $\text{MoO}_3$  and  $\text{KMO}$ ,<sup>8</sup> implying that the quasi-1D electronic property of  $\text{KMO}$  may originate from the quasi-1D zigzag chain of K atoms inserted between  $\text{MoO}_6$  chains. Therefore, XANES at the K  $K$  edge is important to understand the nature of the quasi-1D electronic property. Figure 4 displays the normalized XANES spectra at the K  $K$  edge of  $\text{KMO}$  measured at various angles for the  $b$ - and  $d$ -axis cases. These spectral features are very similar to those observed in other alkali metals and potassium-based compounds.<sup>14,15</sup> Since the electronic state transitions originate from the K  $1s$  core level, only the states with  $p_{xy}$  ( $\sigma^*$  band) and  $p_z$  ( $\pi^*$  band) symmetries centered at the K site are relevant. The inset in Fig. 4 presents magnified near-edge features of  $\theta=0^\circ$ ,  $40^\circ$ , and  $70^\circ$  for both the  $b$ - and  $d$ -axis cases following subtraction of the Gaussian background indicated by the dashed line shown in Fig. 4. A weak preedge feature  $A_3$  (at  $\sim 3608$  eV) was attributed to a *dipole forbidden* K  $1s$  transition to  $4s$  that hybridizes with O  $2p$  states and the sharp spectral features  $B_3$  (at  $\sim 3610$  eV),  $C_3$  (at  $\sim 3612$  eV),  $D_3$  (at  $\sim 3615$  eV), and  $E_3$  ( $\sim 3618$  eV) in the spectra were attributed to transitions from the  $1s$  to the  $4p$  states of K by Teodorescu *et al.*<sup>14</sup> and Marcelli *et al.*<sup>15</sup> The inset in the figure presents the variation of the intensities of these spectral features at various angles for the  $b$ - and  $d$ -axis cases. The changes in the spectral features can be understood by considering that K  $4p_{xy}$  ( $B_3$  and  $C_3$ ) and K  $4p_z$  ( $D_3$  and  $E_3$ ) bands are preferentially probed at large and small incidence angles, respectively, so that along the  $b$  axis, features  $D_3$  and  $E_3$  have maximum intensity at  $\theta=0^\circ$  and  $B_3$  and  $C_3$  have maximum intensity at  $\theta=70^\circ$ . Along the  $d$  axis, the features  $B_3$  and  $C_3$  have similar angle dependence as in the  $b$  axis. However, features  $D_3$  and  $E_3$  do not show a clear trend because their intensities are largest at  $\theta=40^\circ$ , which may be due to the two nonequivalent K sites in the zigzag chain, as

shown in Fig. 1. We have also compared the spectral features with those measured at room temperature above the Peierls transition temperature (not shown here), which shows that peak  $A_3$  does not have significant variation with temperature in both axes, indicating that the dipole-forbidden transition is allowed because of hybridization, not phonon coupling.

The above angle-dependent XANES spectra at the K  $K$  edge of  $\text{KMO}$  show more significant variations in the intensity, general line shape and energy position than those at the O  $K$  and Mo  $L_3$  edge, which suggests that the anisotropy of this compound is strongly related to inserted K atoms. The  $B_3$  and  $C_3$  features associated with K  $4p_{xy}$  orbitals in the K  $K$ -edge XANES spectra have similar angle dependence as those of the O  $K$ -edge XANES spectra, which suggests hybridization between K  $4p$  and O  $2p$  orbitals. Since cation K has a much smaller electronegativity than cation Mo, K atoms are expected to donate electrons and significantly modify charge transfer in the  $\text{MoO}_3$  host system and this effect is anisotropic due to the zigzag chain structure of K ions. The inserted K atoms also give rise to an altered ligand field, which splits K  $4p$  states into two components,  $\sigma^*$  and  $\pi^*$ , and thus alters Mo  $4d$ -O  $2p$  hybridized bands.

In summary, this work employed angle-dependent XANES, which is extremely sensitive to both the geometry and the chemical interlayer of the ionic environment, to demonstrate that significant anisotropy in the electronic properties of  $\text{KMO}$  by comparing the spectral details of the K  $K$  edge with those of the O  $K$  and Mo  $L_3$  edges. The alkali ions K are found to cause asymmetry in the crystal field of Mo through anisotropic hybridization with O ions in the octahedron due to the quasi-1D zigzag chain structure of K ions inserted between host  $\text{MoO}_3$  chains.

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